

Multivariance as a crucial property of microcosm

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Abstract

The conventional method of a generalized geometry construction, based on deduction of all propositions of the geometry from axioms, appears to be imperfect in the sense, that multivariant geometries cannot be constructed by means of this method. Multivariant geometry is such a geometry, where at the point P there are many vectors \mathbf{PP}' , \mathbf{PP}'' , ... which are equivalent to the vector \mathbf{QQ}' at the point Q , but they are not equivalent between themselves. In the conventional (Euclidean) method the equivalence relation is transitive, whereas in a multivariant geometry the equivalence relation is intransitive, in general. It is a reason, why the multivariant geometries cannot be deduced from a system of axioms. The space-time geometry in microcosm is multivariant. Multivariant geometry is a grainy geometry, i.e. the geometry, which is partly continuous and partly discrete. Multivariance is a mathematical method of the graininess description. The graininess (and multivariance) of the space-time geometry generates a multivariant (quantum) motion of particles in microcosm. Besides, the grainy space-time generates some discrimination mechanism, responsible for discrete parameters (mass, charge, spin) of elementary particles. Dynamics of particles appears to be determined completely by properties of the grainy space-time geometry. The quantum principles appear to be needless.

1 Introduction

At first about the term "crucial" with respect to a property of a physical theory. In fifteenth and sixteenth centuries, when the transition from the Aristotelian mechanics to the Newtonian mechanics took place, the crucial concept was "inertia". This concept was absent in the Aristotelian mechanics, but this concept was the new concept, appeared in the Newtonian mechanics. Formally the increase of the order

of dynamic equations of the physical body motion was connected with the concept of inertia. A chariot moving by a horse is a symbol of the Aristotelian mechanics (there is no inertia). A pendulum, whose vibration can be explained only by means of the concept of inertia, is a symbol of the Newtonian mechanics. Introduction of the crucial concept into mechanics lasted longer, than a century. This introduction was accompanied by difficulties and conflicts between the investigators. For instance, conflict between Ptolemaic successors and successors of Copernicus was conditioned by a use of the concept of inertia. According to Ptolemaic conception the planetary system is a great mechanism, which was put in motion by God, whereas according to Copernicus conception the planets move themselves by inertia. Finally, the concept of inertia was so important, that Sir Isaac Newton devoted the first law of mechanics to formulation of this concept, although in reality the first law of mechanics is simply a special case of the second law of mechanics. Appearance of the concept of inertia is conditioned by the transition from earthen mechanics, where the friction force was a dominating reason of dynamics, to the celestial mechanics, where the friction force may be neglected.

At transition from the macroscopic mechanics to mechanics of microcosm a new crucial concept appears. This new concept is called multivariance. When one investigated a passage of electrons through a narrow slit, one discovered, that the electron motion ceases to deterministic (electron diffraction). The electron motion becomes to be multivariant (nondeterministic). Principles of the classical mechanics do not admit a multivariant motion of a free particle. However, the experiment shows, that the motion of small (elementary) particles may be multivariant. Motion is determined by two factors: (1) space-time geometry and (2) laws of dynamics. Thus, there are two possibilities: either the space-time geometry is multivariant, or dynamics in microcosm is multivariant (it may be also, that both geometry and dynamics are multivariant). In the thirtieth of the twentieth century, when the electron diffraction was discovered, the multivariant geometry was not known. Nobody could imagine, that the space-time geometry may be multivariant. (Note, that nondeterministic geometries were known. But in reality, there were stochastic structures, given on a geometry, whereas the geometry in itself was deterministic and single-variant). As a result the multivariance has been ascribed to dynamics. This multivariant dynamics is known as the quantum mechanics. Note that appearance of the multivariant geometry in dynamics is not identical to quantum dynamics. The quantum dynamics is a special case of multivariant dynamics. The multivariant dynamics contains the quantum dynamics and something else, which cannot be reduced to the quantum dynamics. This "something else" is of interest.

It appeared that the geometry may be multivariant [1]. Multivariance of geometry means as follows. Let $\mathbf{P}_0\mathbf{P}_1$ be a vector at the point P_0 . Let at the point Q_0 there be many vectors $\mathbf{Q}_0\mathbf{Q}_1, \mathbf{Q}_0\mathbf{Q}_2, \dots$, which are equivalent (equal) to the vector $\mathbf{P}_0\mathbf{P}_1$ at the point P_0 , but vectors $\mathbf{Q}_0\mathbf{Q}_1, \mathbf{Q}_0\mathbf{Q}_2, \dots$ are not equivalent between themselves. If such a situation takes place in geometry, then such a geometry is multivariant.

If at any point Q_0 there is one and only one vector $\mathbf{Q}_0\mathbf{Q}_1$, which is equivalent to the vector $\mathbf{P}_0\mathbf{P}_1$ at the point P_0 , such a geometry is called the single-variant

geometry.

In general, multivariance and single-variance of the geometry are considered with respect to some pair of points P_0, Q_0 . It is possible such a situation, when the geometry is multivariant with respect to some pairs of points, and it is single-variant with respect to another pairs of points. If the geometry is multivariant with respect to at least one pair of points, such a geometry will be qualified as multivariant. In the multivariant space-time geometry the particle dynamics appears to be multivariant, even if this dynamics acts in accordance with conventional principles of the classical dynamics.

Note, that the equivalence relation is supposed to be transitive in all mathematical models. i.e. in all logical constructions, which can be deduced from a system of axioms by means of the rules of the formal logic. The equivalence relation is transitive by definition, if for any objects (for instance, vectors $\mathbf{P}_0\mathbf{P}_1, \mathbf{Q}_0\mathbf{Q}_1, \mathbf{Q}_0\mathbf{Q}_2$) it follows from the relations $\mathbf{P}_0\mathbf{P}_1 \text{eqv} \mathbf{Q}_0\mathbf{Q}_1$ and $\mathbf{P}_0\mathbf{P}_1 \text{eqv} \mathbf{Q}_0\mathbf{Q}_2$, that $\mathbf{Q}_0\mathbf{Q}_1 \text{eqv} \mathbf{Q}_0\mathbf{Q}_2$. Here designation "eqv" means relation of equivalence. Comparison of definition of multivariance ($\mathbf{P}_0\mathbf{P}_1 \text{eqv} \mathbf{Q}_0\mathbf{Q}_1 \wedge \mathbf{P}_0\mathbf{P}_1 \text{eqv} \mathbf{Q}_0\mathbf{Q}_2$, but $\mathbf{Q}_0\mathbf{Q}_1 \overline{\text{eqv}} \mathbf{Q}_0\mathbf{Q}_2$) with the definition of transitivity shows that the equivalence relation in the multivariant geometry cannot be always transitive.

However, does the multivariant geometry (T-geometry) exist? If yes, then how can one construct a multivariant geometry?

Let us consider the proper Euclidean geometry and define equivalence of two vectors $\mathbf{P}_0\mathbf{P}_1$ and $\mathbf{Q}_0\mathbf{Q}_1$ as follows. Vectors $\mathbf{P}_0\mathbf{P}_1$ and $\mathbf{Q}_0\mathbf{Q}_1$ are equivalent ($\mathbf{P}_0\mathbf{P}_1 \text{eqv} \mathbf{Q}_0\mathbf{Q}_1$), if vectors $\mathbf{P}_0\mathbf{P}_1$ and $\mathbf{Q}_0\mathbf{Q}_1$ are in parallel ($\mathbf{P}_0\mathbf{P}_1 \uparrow\uparrow \mathbf{Q}_0\mathbf{Q}_1$) and their lengths $|\mathbf{P}_0\mathbf{P}_1|$ and $|\mathbf{Q}_0\mathbf{Q}_1|$ are equal. Mathematically these two conditions are written in the form

$$(\mathbf{P}_0\mathbf{P}_1 \uparrow\uparrow \mathbf{Q}_0\mathbf{Q}_1) : \quad (\mathbf{P}_0\mathbf{P}_1 \cdot \mathbf{Q}_0\mathbf{Q}_1) = |\mathbf{P}_0\mathbf{P}_1| \cdot |\mathbf{Q}_0\mathbf{Q}_1| \quad (1.1)$$

$$|\mathbf{P}_0\mathbf{P}_1| = |\mathbf{Q}_0\mathbf{Q}_1|, \quad |\mathbf{P}_0\mathbf{P}_1| = \sqrt{2\sigma(P_0, P_1)} \quad (1.2)$$

where $(\mathbf{P}_0\mathbf{P}_1 \cdot \mathbf{Q}_0\mathbf{Q}_1)$ is the scalar product of two vectors, defined by the relation

$$(\mathbf{P}_0\mathbf{P}_1 \cdot \mathbf{Q}_0\mathbf{Q}_1) = \sigma(P_0, Q_1) + \sigma(P_1, Q_0) - \sigma(P_0, Q_0) - \sigma(P_1, Q_1) \quad (1.3)$$

Here σ is the world function of the proper Euclidean space, which is defined via the Euclidean distance $\rho(P, Q)$ between the points P, Q by means of the relation

$$\sigma(P, Q) = \frac{1}{2}\rho^2(P, Q) \quad (1.4)$$

The length $|\mathbf{PQ}|$ of vector \mathbf{PQ} is defined by the relation

$$|\mathbf{PQ}| = \rho(P, Q) = \sqrt{2\sigma(P, Q)} \quad (1.5)$$

Using relations (1.1) - (1.5), one can write the equivalence condition in the form $\mathbf{P}_0\mathbf{P}_1 \text{eqv} \mathbf{Q}_0\mathbf{Q}_1$:

$$\sigma(P_0, Q_1) + \sigma(P_1, Q_0) - \sigma(P_0, Q_0) - \sigma(P_1, Q_1) = \sigma(P_0, P_1) \quad (1.6)$$

$$\wedge \sigma(P_0, P_1) = \sigma(Q_0, Q_1) \quad (1.7)$$

The definition of equivalence (1.6), (1.7) is a satisfactory geometrical definition, because it does not contain a reference to the dimension of the space and to the coordinate system. It contains only points P_0, P_1, Q_0, Q_1 , determining vectors $\mathbf{P}_0\mathbf{P}_1$ and $\mathbf{Q}_0\mathbf{Q}_1$ and distances (world functions) between these points. The definition of equivalence (1.6), (1.7) coincides with the conventional definition of two vectors equivalence in the proper Euclidean geometry. If one fixes points P_0, P_1, Q_0 in the relations (1.6), (1.7) and solve them with respect to the point Q_1 , one finds that these equations always have one and only one solution. This statement follows from the properties of the world function of the proper Euclidean geometry. It means that the proper Euclidean geometry is single-variant with respect any pairs of its points. It means also, that the equivalence relation is transitive in the proper Euclidean geometry.

Any geometry is a set (in general, continual one) of propositions. The proper Euclidean geometry may be axiomatized, i.e. all propositions of the proper Euclidean geometry may be deduced from a finite set of propositions (axioms) by means of the rules of formal logic. The system of axioms is consistent [2]. This fact is in accordance with the transitivity of the equivalence relation in the proper Euclidean geometry.

On the other hand, all propositions of the proper Euclidean geometry may be expressed in terms of the world function [1]. Let us represent all propositions of the proper Euclidean geometry in terms of the Euclidean world function σ_E and replace the Euclidean world function σ_E by some another world function σ , satisfying the constraints

$$\sigma : \quad \Omega \times \Omega \rightarrow \mathbb{R}, \quad \sigma(P, Q) = \sigma(Q, P), \quad \sigma(P, P) = 0, \quad \forall P, Q \in \Omega \quad (1.8)$$

where Ω is the set of all points, where the geometry is given. We obtain the set of all propositions of the geometry \mathcal{G} , described by the world function σ . Such a replacement is a deformation of the proper Euclidean geometry. Thus, it is possible to construct the "metric" geometry, which contains all propositions of the proper Euclidean geometry. I shall not use the term "metric geometry" for the deformed geometry \mathcal{G} , because the geometry \mathcal{G} is free of the constraint (the triangle axiom), which is imposed on the metric geometry

$$\rho(P, R) + \rho(R, Q) \geq \rho(P, Q), \quad \forall P, Q, R \in \Omega \quad (1.9)$$

The triangle axiom (1.9) is imposed, in order to conserve one-dimensional character of shortest (straight) in the metric geometry. Indeed, in the proper Euclidean geometry the set

$$\mathcal{E}\mathcal{L}_{P_1, P_2, Q} = \{R | \rho(P_1, R) + \rho(R, P_2) = \rho(P_1, Q) + \rho(Q, P_2)\} \quad (1.10)$$

is an ellipsoid with focuses at the points P_1, P_2 and the point Q on the surface of the ellipsoid. If the point Q tends to the focus P_2 , the ellipsoid degenerates in segment

$$\mathcal{T}_{[P_1, P_2]} = \{R | \rho(P_1, R) + \rho(R, P_2) = \rho(P_1, P_2)\} \quad (1.11)$$

of the straight line, passing through the points P_1 and P_2 . In the proper Euclidean geometry the ellipsoid degenerates into one-dimensional segment of straight. However, in the arbitrary metric geometry, given on n -dimensional manifold, the equation

$$\mathcal{S} : \quad \Phi(R) = 0, \quad \Phi(R) \equiv \rho(P_1, R) + \rho(R, P_2) - \rho(P_1, P_2) \quad (1.12)$$

determines, in general, $(n - 1)$ -dimensional closed surface \mathcal{S} . The points R , satisfying the condition $\Phi(R) > 0$ are external points, which are placed outside the closed surface \mathcal{S} . The points R , satisfying the condition $\Phi(R) < 0$ are internal points, which are placed inside \mathcal{S} . If the condition (1.9) is satisfied, it means, that the closed surface \mathcal{S} has no internal points. In this case the segment (1.11) has no internal points, i.e. it is one-dimensional.

In the deformed geometry \mathcal{G} the solution of equations (1.6), (1.7) for the point Q_1 at fixed points P_0, P_1, Q_0 does not always exist. If it exists, it is not always unique. In other words, the deformed geometry \mathcal{G} is multivariant, in general. In the same time any proposition of the proper Euclidean geometry exists in the deformed geometry \mathcal{G} , although it may have another sense, than the sense, which this proposition has in the proper Euclidean geometry. Nevertheless, this proposition is the same proposition, formulated in different geometries.

Some conventional propositions of the proper Euclidean geometry contain references to the dimension and to the coordinate system, i.e. to the method of the geometry description. In the conventional (vector) presentation of the Euclidean geometry, its dimension is considered to be a property of the geometry in itself, although there are geometries, where the dimension cannot be introduced, because for introduction of the dimension the world function must satisfy a series of constraints, which are very restrictive. In reality, the dimension of geometry is the dimension of the coordinate system (the number of coordinates), which is used for the geometry description. The manifold and its dimension is the conventional method of the geometry description [3], and one cannot separate this method from the geometry in itself, until one has not alternative method of the geometry construction.

The deformed geometry \mathcal{G} is a multivariant geometry, which cannot be axiomatized, in general. It means that the geometry \mathcal{G} is an example of a theory, which cannot be considered as a conventional mathematical model, constructed by means of the formal logic on the basis of some axiomatics. Returning to the multivariance, discovered in the motion of electrons, one may state, that the problem of multivariance of the electron motion can be solved on account of a multivariant geometry. The multivariant space-time geometry looks more reasonable, than the multivariant dynamics in the single-variant space-time geometry. Indeed, to obtain multivariant dynamics one is forced to replace principles of classical dynamics by quantum principles, which looks rather artificial. In the same time the multivariance is a natural property of the physical geometry (i.e. of the geometry, described completely by the world function). It depends on the form of the world function σ in the equivalence relations (1.6), (1.7), whether or not the geometry \mathcal{G} is multivariant. Changing the world function of the space-time, one can change a character of the multivari-

ance. One can choose such a world function of the space-time, that the conventional classical description of a particle motion coincide with the description of quantum mechanics [4].

Let the space-time geometry be described by the world function

$$\sigma_d = \sigma_M + d \operatorname{sgn}(\sigma_M), \quad d \equiv \lambda_0^2 = \frac{\hbar}{2bc} = \text{const} \quad (1.13)$$

$$\operatorname{sgn}(x) = \begin{cases} 1 & \text{if } x > 0 \\ 0 & \text{if } x = 0 \\ -1 & \text{if } x < 0 \end{cases}$$

where σ_M is the world function of the Minkowski space-time, \hbar is the quantum constant, c is the speed of the light and b is some universal constant. Then the world chain, consisting of points $\dots P_0, P_1, \dots P_k, \dots$ satisfying the relations

$$\mathbf{P}_k \mathbf{P}_{k+1} \text{eqv} \mathbf{P}_{k+1} \mathbf{P}_{k+2}, \quad k = \dots 0, 1, \dots \quad (1.14)$$

describes the motion of a free particle. It appears, that the motion is multivariant (stochastic) in the space-time with the world function (1.13). Statistical description of these multivariant chains coincides with the quantum description in terms of the Schrödinger equation [4].

Besides, the space-time (1.13) appears to be discrete, because in this space-time there are no vectors \mathbf{PQ} of the length $|\mathbf{PQ}|^4 \in (0, \lambda_0^4)$. Discreteness of the space-time seems to be very surprising, because the space-time is given on the manifold of Minkowski. Conventionally the discrete space is associated with a grid. The discrete space-time on the continuous manifold seems to be impossible. This example shows, that a physical geometry and a continuous manifold, where the geometry is given, are quite different things. Manifold and its dimension are only attributes of the vector representation of the Euclidean geometry (i.e. of the description method) [3], whereas the discreteness of a geometry is an attribute of the geometry in itself.

In any mathematical model the equivalence relation is transitive. This property of the mathematical model provides definiteness (single-variance) for all conclusions, made on the basis of such a mathematical model. If the equivalence relation is intransitive, the conclusions, made on the basis of such a model cease to be definite. They becomes multivariant. The logical construction with the intransitive equivalence relation and, hence, with multivariant conclusions is not considered to be a mathematical model, because it is useless, and one cannot make a definite prediction on the basis of such a model. Besides, such a model cannot be axiomatized, because the axiomatization supposes a single-variance of conclusions.

I shall refer to models with multivariant (indefinite) predictions as intransitive models, or multivariant models. Multivariant models appear automatically, as soon as they use a multivariant space-time geometry. As far as models of physical phenomena may not ignore space-time geometry, and the space-time geometry may be multivariant, one cannot avoid a use of multivariant models of physical phenomena.

Fortunately, a multivariant model can be reduced to a single-variant model, provided one unites the set of many conclusions, which follows from one statement into

one conclusion. In other words, one considers the set of different objects as a statistical ensemble. One may work with the statistical ensemble, considering it as a single object. Then the multivariant model may cease to be multivariant. It turns into a single-variant (transitive) model provided, that its objects be statistical ensembles of original objects. Such a procedure is known as the statistical description, which deals with statistically averaged objects. Prediction of the model about statistically averaged objects (statistical ensembles), which now are objects of the model, may appear to be single-variant, if the statistical description is produced properly. In other words, a statistical description, produced properly, transforms a multivariant model into a single-variant mathematical model.

Procedure of the statistical description is well known. It is used in different branches of theoretical physics. However, sometimes one obtains the single-variant mathematical model, dealing with statistically averaged objects, without knowing that the model deals with statistically averaged objects. For instance, the gas dynamics model deals with gas particles. Motion of gas particle describes the mean motion of gas molecules. However, the gas dynamics equations (as dynamic equations of the continuous medium) were deduced, before it became known, that the gas consists of molecules. Besides, there is more detailed statistical description of the gas molecule motion, based on the gas kinetic theory (Boltzman equation).

The quantum mechanics is a statistical description of the multivariant particle motion, which is conditioned directly by the multivariant space-time geometry (1.13). However, the quantum mechanics is not considered conventionally as a statistical description of a multivariant particle motion. One considers the quantum mechanics as a corollary of special quantum principles of dynamics, which are introduced axiomatically. In this form the quantum mechanics describes very successfully physical phenomena of atomic physics. Formal technique of quantum mechanics is rather simple and comfortable. Many investigators like the formalism of quantum mechanics, and they argue against the quantum mechanics as a statistical description of multivariantly moving particles.

Something like this, one had more, than hundred years ago with the thermal phenomena. The heat was explained as a special heat liquid (thermogen), whose properties are described by the laws of thermodynamics. The axiomatic thermodynamics explained very well all thermal phenomena. Attempts of interpretation of the heat as a chaotic molecular motion met objections of many investigators, who did not believe in existence of molecules. The heat as a chaotic molecular motion had been accepted, when it became clear that the thermal fluctuations cannot be explained by the axiomatic thermodynamics. The thermal fluctuation can be explained only by the supposition that the heat is a chaotic molecular motion. However, the axiomatic thermodynamics is much simpler, than the statistical theory of the heat. It is used now in the theory of continuous medium and other applications.

Situation with the quantum mechanics looks as follows. In general, the quantum mechanics may be deduced as a result of statistical description of the multivariant motion of particles, conditioned by the multivariant space-time geometry of the form (1.13). However, the quantum mechanics had been formulated in the beginning of

the twentieth century as an axiomatic conception. The multivariant space-time geometry was not known then. Now most of investigators do not see a necessity of introducing the concept of multivariant geometry. The fact, that introduction of quantum principles is a corollary of our imperfect knowledge of geometry, does not disturb them. They believe that the relativistic quantum theory and the theory of elementary particles can be constructed on the basis of unification of the quantum principles and principles of relativity.

Strategy of further investigations of the microcosm depends essentially on the approach to the multivariant space-time geometry. If we believe, that the multivariant space-time geometry is impossible, and quantum principles reflect the nature of the microcosm, we are forced to use the investigation strategy, which has been used at the construction of the nonrelativistic quantum mechanics. The quantum mechanics has been constructed by the cut-and-try method. The same method is used for further investigation of the microcosm. Besides, the quantum principles supposes, that all physical objects and all physical fields are to be quantized. In particular, one should quantized the gravitational and electromagnetic fields.

On the contrary, if one takes, that the quantum effects are a corollary of the multivariant geometry, one should not quantized the electromagnetic and gravitational fields, because these fields describe the space-time geometry. Besides, the dynamic equations of the electromagnetic field and those of the gravitational field do not contain the quantum constant. This fact manifests a distinction of dynamic equations of these fields from the Schrödinger and Dirac equations. From the logical viewpoint the approach, based on a use of the multivariant space-time geometry, is more consistent also. Indeed, why would one use only single-variant space-time geometries, which form only a small part of all possible space-time geometries? When it appears, that the single-variant space-time cannot explain the multivariant motion of free particles, one is forced to introduce enigmatic quantum principles to explain quantum effects, which are a manifestation of multivariance. In general, why is one to ignore the property of multivariance, which is observed in experiments on the electron diffraction?

Note that according to the definition of equivalence (1.6), (1.7), there exists the zero-variance, when the equations (1.6), (1.7) have no solutions. If the multivariance may be reduced to single-variance of statistically averaged objects by means of statistical description, the zero-variance of the space-time geometry cannot be described by a single-variant mathematical model. The zero-variance means discrimination, when some variants of the particle motion are discriminated. For instance, the space-time geometry (1.13) discriminates the particles of small masses, because in the multivariant space-time geometry the masses of particles are geometrized, and the particle mass m is connected with the lengths $|\mathbf{P}_k \mathbf{P}_{k+1}|$ of the vectors of the world chain by the relation

$$m = b |\mathbf{P}_k \mathbf{P}_{k+1}| \quad (1.15)$$

where b is the universal constant, which enters in the expression (1.13) for the elementary length λ_0 .

The fact, that masses of elementary particles, their electric charges and their

internal angular moments (spin) are discrete quantities (but not all possible ones) is a result of some discrimination mechanism connected with the possible zero-variance of the space-time geometry. The values of electric charge and those of spin are multiple to quantities e and \hbar respectively. This fact is postulated in the framework of quantum mechanics. Discreteness of the elementary particles masses is postulated also. However, the values of masses are taken from experiment, and theorists dream to deduce the receipt of determination of the mass values, considering this receipt as a great achievement of the elementary particle theory. However, the quantum principles do not admit to determine discrete values of the elementary particles masses. These discrete values of masses (as well as the values of the electric charge and spin) should be determined by some discrimination mechanism, which is conditioned by the multivariant (zero-variant) space-time geometry. Such a possibility must be investigated, because, being a corollary of a statistical description, the quantum principles do not admit one to determine such a discrimination mechanism.

Investigation of the space-time geometry admits one to set the question, what elementary particles may exist at given space-time geometry. To determine the proper space-time geometry, one may vary the values of the world function (1.13) in the interval, where $\sigma \in (-\lambda_0^2, \lambda_0^2)$. Variation of the form of the world function σ for the values of argument σ_M in the interval, where $\sigma \in (-\lambda_0^2, \lambda_0^2)$ does not influence on the Schrödinger equation, generated by the multivariant geometry (1.13). In the conventional approach, when only single-variant space-time geometry is considered, the question on geometrical justification of the elementary particles existence cannot be put at all. This question is set only on the level of dynamics, where one has no discrimination mechanism. In the multivariant space-time geometry one can consider the question of the limited divisibility of geometrical objects [5]. In the single-variant geometry such a question cannot be put, because in such a geometry the unlimited divisibility takes place.

2 Why do the most scientists ignore concept of multivariant space-time geometry and concept of multivariance?

This question is not a scientific question. This is a social near-scientific question. I do not know the answer for this question. But this question is very important for further development of the microcosm physics, because it admits one to choose an effective investigation strategy. I try to consider different versions of the answer. In reality, one separates this global question into a series of more special questions and tries to answer some of them.

It is impossible to find a defect in construction of T-geometry in itself. It is too simple, in order one could find a mistake or a defect in its construction. There are three points in the method of construction of T-geometry:

- (1) T-geometry is a physical (metric) geometry, which is described completely

by the world function and only by the world function.

(2) Method of construction of geometrical objects and of the T-geometry propositions is the same for all T-geometries, i.e. the formula of description in terms of the world function is the same in all T-geometries.

(3) The proper Euclidean geometry is a mathematical (axiomatized) geometry and a physical geometry simultaneously. There is a theorem of the Euclidean geometry, which states, that the proper Euclidean geometry may be described in terms of the world function and only in terms of the world function [1].

The point (2) follows from the point (1). Indeed, let the geometrical object be described by the formula a_1 in a physical geometry \mathcal{G}_1 , and the same object be described by the formula a_2 in other physical geometry \mathcal{G}_2 . If formulas a_1 and a_2 are different, it means that the geometries \mathcal{G}_1 and \mathcal{G}_2 distinguish not only by their world functions. There is some quantity, which is different for \mathcal{G}_1 and \mathcal{G}_2 , and this circumstance disagrees with the point (1).

It follows from the point (3), that all propositions of the proper Euclidean geometry can be deduced from Euclidean axioms and presented in terms of the proper Euclidean world function σ_E . Replacing σ_E in all propositions of the proper Euclidean geometry by the world function σ of the physical geometry \mathcal{G} , one obtains all propositions of the geometry \mathcal{G} and, hence, the physical geometry \mathcal{G} in itself. The point (3) admits one to construct any physical geometry, basing on our knowledge of the proper Euclidean geometry.

The non-Euclidean method of the physical geometry construction (the deformation principle) [6] is simpler, than the conventional Euclidean method of the geometry construction, because it does not need a proof of theorems and a test of the axioms consistency. One may say, that the conventional method takes from Euclid the intermediate product (method of the geometry construction), whereas the non-Euclidean method takes from Euclid his final product (the Euclidean geometry in itself). The intermediate product needs a further work (proof of theorems), whereas the final product does not need further work, because all necessary theorems are supposed to be proved at the stage of the proper Euclidean geometry construction.

Thus, the deformation principle has not difficulties of the Euclidean method. Besides, it admits one to construct multivariant geometries. However, the most mathematicians do not accept the deformation principle. For instance, the author of this paper submitted a report on the construction of T-geometry to a geometro-topological seminar in the Moscow Lomonosov University. The secretary of seminar looked through the presented paper and said something like that: "How strange geometry! There are no theorems! Only definitions! I think, that such a geometry is not interesting for participants of our seminar." The secretary was quite right. The main activity of geometro-topologists is a formulation of theorems and a proof of them. Such an activity cannot be applied in the geometry, constructed by means of the deformation principle.

The secretary of another geometro-topological seminar investigated papers, submitted for reading of my report. The report was devoted to construction of T-

geometry. Submitted documents contained, in particular, the paper [7]. Investigation of submitted papers lasted almost a year. It was finished with the decision: "Participants of the seminar are not ready to hear the report." Such a decision means, that the participants of the seminar are not able to argue anything against the T-geometry, but nevertheless, they cannot accept it. Another examples of negative relation to the T-geometry construction one can find in [7].

I must note, that there are mathematicians, whose relation to the T-geometry construction was well-minded. They were participants of the seminar on "geometry as a whole" in the Moscow Lomonosov University. Reports on the T-geometry construction were read at this seminar several times. However, participants of this seminar were not geometro-topologists.

The geometro-topologists construct generalized geometries on the basis of a topological space and corresponding axiomatics, and the negative relation to T-geometry may be interpreted in the sense, that accepting T-geometry, one depreciates papers, based on the conventional (topological) approach to the construction of generalized geometries. However, I should not like to interpret the negative reaction of topologists in such a way. I should prefer to understand objective reasons of negative relation to the T-geometry.

Expression of Euclidean propositions in terms of the world function supposes another set \mathcal{A}_d of primary axioms, than the set \mathcal{A}_c of primary axioms, which are used usually. For instance, the set \mathcal{A}_c contains the axiom: "The straight has no width." The system of primary axioms \mathcal{A}_d does not contain this statement. The statement "the straight has no width." is valid (for the proper Euclidean geometry) in the system of axioms \mathcal{A}_d , but it is a secondary statement in \mathcal{A}_d . It is a result of the axiomatics \mathcal{A}_d and of definition of the straight. The definition of the straight $\mathcal{T}_{P_0P_1}$, passing through the points P_0, P_1 has the form

$$\mathcal{T}_{P_0;P_0P_1} = \mathcal{T}_{P_0P_1} = \{R|\mathbf{P}_0\mathbf{R} \parallel \mathbf{P}_0\mathbf{P}_1\} \quad (2.1)$$

where the relation of collinearity $\mathbf{P}_0\mathbf{P}_1 \parallel \mathbf{P}_0\mathbf{R}$ is defined by the relation

$$\mathbf{P}_0\mathbf{R} \parallel \mathbf{P}_0\mathbf{P}_1: \quad (\mathbf{P}_0\mathbf{R} \cdot \mathbf{P}_0\mathbf{P}_1)^2 = |\mathbf{P}_0\mathbf{P}_1|^2 |\mathbf{P}_0\mathbf{R}|^2 \quad (2.2)$$

Here the scalar product $(\mathbf{P}_0\mathbf{P}_1 \cdot \mathbf{P}_0\mathbf{R})$ is defined by the relation (1.3). In general, one equation (2.2) defines a $(n - 1)$ -dimensional surface on the n -dimensional manifold (but not a one-dimensional straight). The statement, that the point set (2.1), (2.2) is a one-dimensional straight, which has no width, follows from the properties of the Euclidean world function. This property may not take place for other world function.

Note that the point set

$$\mathcal{T}_{P_0;Q_0Q_1} = \{R|\mathbf{P}_0\mathbf{R} \parallel \mathbf{Q}_0\mathbf{Q}_1\} \quad (2.3)$$

$$\mathbf{P}_0\mathbf{R} \parallel \mathbf{Q}_0\mathbf{Q}_1: \quad (\mathbf{P}_0\mathbf{R} \cdot \mathbf{Q}_0\mathbf{Q}_1)^2 = |\mathbf{Q}_0\mathbf{Q}_1|^2 |\mathbf{P}_0\mathbf{R}|^2 \quad (2.4)$$

is also a $(n - 1)$ -dimensional surface on the n -dimensional manifold. In the proper Euclidean geometry the set $\mathcal{T}_{P_0;Q_0Q_1}$ degenerates into the straight line, passing through the point P_0 in parallel with the vector $\mathbf{Q}_0\mathbf{Q}_1$ at the point Q_0 .

In the n -dimensional Riemannian geometry the $(n - 1)$ -dimensional point set (2.1), (2.2) also degenerates into one-dimensional geodesic, passing to the point P_0 in parallel with the vector $\mathbf{P}_0\mathbf{P}_1$. This degeneration is conditioned by the fact, that the Riemannian space may be considered as a metric space with the metric, satisfying the triangle axiom. However, the point set (2.3), (2.4) does not degenerate, in general, into one-dimensional curve (geodesic). In the n -dimensional Riemannian geometry the point set (2.3), (2.4) remains to be a $(n - 1)$ -dimensional surface, as well as in any T-geometry (except for the proper Euclidean geometry).

This fact means, that the Riemannian geometry is a multivariant physical geometry, although some sort of straights ($\mathcal{T}_{P_0;P_0P_1} = \mathcal{T}_{P_0P_1}$) is single-variant (one-dimensional). On the other hand, the Riemannian geometry is constructed usually as a single-variant geometry, and existence of geometrical objects (2.3), (2.4) is incompatible with the axiom "geodesic has no width". Geodesic is defined as a curve of minimal (extremal) length. In turn the curve is defined as a continuous mapping

$$[0, 1] \rightarrow \Omega$$

which cannot be formulated in terms of the world function only, because it contains a reference to a manifold. To remove the disagreement between the multivariance of definition (2.3), (2.4) and axiom "geodesic has no width", one declared, that there is no fernparallelism in the Riemannian geometry, i.e. parallelism of remote vectors is not determined. At such a constraint the geometrical objects (2.3), (2.4) are not defined, and hence, they do not exist.

However, the removal of the fernparallelism does not eliminate inconsistency of the Riemannian geometry, it eliminates only one of corollaries of this inconsistency. There may be another (unknown) corollaries of this inconsistency. One can eliminate these corollaries only removing their reason (axiom, that the geodesic has no width). It means that one should accept the definition of the straight (geodesic) in the form (2.3), (2.4), i.e. the idea of multivariance should be accepted.

Inconsistency of a conception manifests itself only, if one solves a problem by different correct methods and obtains different results. However, rare scientists investigate a complicate problem by several different methods and compare the obtained results.

The outstanding topologist G.Perelman had proved the Poincarè conjecture [8, 9, 10]. In 2006 he was awarded with the Fields medal. However, he declined to accept the award. He is the only person ever to refuse the award. Besides, he declined to publish his papers in a peer review journal, that was necessary for accepting a prize of a million dollars. His behavior looked strange and unexpected for mathematical society. Alexander Abramov [11], who knew Perelman personally very well, describes his style of work as follows. Perelman considered several versions of solution of the problem and chose the best one. Such a rare style of investigation is the best one for discovery of inconsistencies in the Riemannian geometry. Apparently, after publication of his papers in Archives Perelman has discovered, that the conventional (topological) approach to the Riemannian geometry is inconsistent (maybe, the paper [7], appeared in March 2005, gave a motive for his investigation).

But G.Perelman is a topologist and his papers on the Poincarè conjecture are based on the Riemannian geometry. If the Riemannian geometry is inconsistent, his own papers become questionable, even if all his considerations are valid.

His further behavior is conditioned by his scientific scrupulosity. He could not withdraw his papers from Archives, where they were published (it is prohibited by the rules of Archives). But he could decline publication of his papers in the peer review mathematical journals. He could not accept the Fields medal, because some time later his papers may be declared to be questionable. He should publish the fact, that he discovered inconsistency of the Riemannian geometry. But such a paper would be a dissident paper. Anybody, who have written a dissident paper, knows very well, how difficult to publish a dissident paper. Discussing with colleagues a possible inconsistency of the Riemannian geometry, G.Perelman did not meet mutual understanding by colleagues. As a result of such discussions he left the Institute, where he worked. The charge of his colleagues in scientific dishonesty is also a result of these discussions.

I did not know G. Perelman personally, and my description of his dignified behavior is only a hypothesis. But it is a very reasonable hypothesis, which explains freely all facts by the scientific scrupulosity of G. Perelman and by his capacity of investigation work. My estimation of the Perelman's activity distinguishes from position of other scientists, because I know definitely, that the Riemannian geometry is inconsistent, especially in that its part, which concerns topology, whereas other scientists cannot admit an inconsistency of the Riemannian geometry. Topology in the Riemannian geometry, as well as in other physical geometries, is determined completely by the world function. The topology may not be given independently, because in this case one risks to obtain inconsistency.

Construction of multivariant geometry is connected with a replacement of the formal logic by the "Euclidean logic" [12], when rules of the formal logic are substituted by the rules of construction of the Euclidean geometry propositions in terms of the world function. The transition from the formal logic to the "Euclidean logic" is a transition from one system of axioms to another equivalent (for Euclidean geometry) system of axioms. Such a transition is used very uncommon in the practice of mathematical investigations. Although possibility of such transition is accepted, but in practice the transformation of the system of axioms, connected with such a transition, is used insufficiently. In application of any axiomatics there are logical stereotypes, when a chain of logical conclusions is replaced by one statement. Such stereotypes depend on the used axiomatics, and they are changed at a change of axiomatics. At a replacement of the axiomatics the logical stereotypes are to be analyzed and replaced by new logical stereotypes. Unfortunately, the practice of work with logical stereotypes is insufficient. As a result the old logical stereotypes disturb the perception of new axiomatics.

Let me adduce a simple example. In the conventional approach to geometry, based on the vector representation, the discrete geometry cannot be given on a continual set of points (on manifold). It can be given only on a discrete set of point of the type of a grid. On the other hand, by definition, the discrete geometry is

such a geometry, where there are no close points. In the approach, based on the principle of deformation, the distance between points is determined by the world function and only by the world function. It is of no importance, where the world function is given (on a grid, or on a continuous manifold). If the world function is given on a manifold in such a form, that there are no values of the world function σ in the intervals $(-a, 0)$ and $(0, a)$, $a > 0$, then in the geometry there are no close points, and the geometry is discrete, even it is given on a continuous manifold.

The statement (st): "the discrete geometry cannot be given on a manifold" is a logical stereotype of the approach, based on the vector representation of the geometry. This stereotype consists of two statements: (1) definition: the discrete geometry does not contain close points, (2) axiom: the continuous geometry is given on a manifold. Although the statement (st) does not follow strictly from statements (1) and (2) does not follow strictly, because it is not known, where the discrete geometry is given. Nevertheless, because of insufficient development of the discrete geometry one concludes, that the discrete geometry cannot be given on a manifold, as far as on a manifold the continuous geometry is given.

I could not overcome the stereotype (st) during almost fifteen years, when I developed T-geometry. I could not overcome this stereotype, although during fifteen years I dealt with the discrete geometry, which was described by the world function (4.1), given on a continuous manifold. I could not overcome the stereotype, although I developed the world function formalism without any problems. I could not overcome the stereotype, although its essence lies on the surface of the phenomenon. This stereotype is not a unique one. I met another stereotypes at other scientists. I think, that such stereotypes do not admit one to accept idea of deformation principle. In turn the difficulties with overcoming of such stereotypes are connected with the circumstance, that the transition from one axiomatics to another equivalent axiomatics is used very rare in practice. The training of mathematicians for such transitions is too small.

3 Multivariance and dimension

Returning to the T-geometry, I should like to manifest, that the concept of dimension may have different meaning at the conventional approach to geometry and at the approach based on the deformation principle. I shall show, that the dimension of geometry and the dimension of the manifold, where the geometry is given, are different things. The dimension of the manifold $n_{\mathcal{M}}$ and dimension $n_{\mathcal{G}}$ of the geometry are different concepts, which coincide for the proper Euclidean geometry. However, in other physical geometries the two quantities do not coincide, in general.

Let us consider very simple example of the two-dimensional proper Euclidean geometry \mathcal{G}_{E} , given on the two-dimensional manifold. The world function has the form

$$\sigma_{\text{E}}(P_1, P_2) = \sigma_{\text{E}}(\mathbf{x}, \mathbf{y}) = \frac{1}{2} \left((x^1 - y^1)^2 + (x^2 - y^2)^2 \right), \quad \sigma_{\text{E}} \geq 0 \quad (3.1)$$

where the points P_0, P_1, P_2 are three points, whose coordinates in the Cartesian coordinate system are

$$P_0 = \{0, 0\} \quad P_1 = \{x^1, x^2\}, \quad P_2 = \{y^1, y^2\} \quad (3.2)$$

Vectors $\mathbf{P}_0\mathbf{P}_1$ and $\mathbf{P}_0\mathbf{P}_2$ have the Cartesian coordinates

$$\mathbf{P}_0\mathbf{P}_1 = \mathbf{x} = \{x^1, x^2\}, \quad \mathbf{P}_0\mathbf{P}_2 = \mathbf{y} = \{y^1, y^2\} \quad (3.3)$$

Besides, one considers a deformed physical geometry \mathcal{G}_d , described by the world function

$$\sigma_d(P_1, P_2) = \sigma_E(P_1, P_2) + d(\sigma_E(P_1, P_2)) \quad (3.4)$$

where

$$d(\sigma_E) = \begin{cases} -\lambda_0^2 & \text{if } \sigma_E > \sigma_0 \\ -\lambda_0^2 \frac{\sigma_E}{\sigma_0} & \text{if } 0 \leq \sigma_E \leq \sigma_0 \end{cases}, \quad \lambda_0^2 \geq \sigma_0 \geq 0, \quad \lambda_0, \sigma_0 = \text{const} \quad (3.5)$$

Here λ_0 is some elementary length, which is characteristic for the distorted geometry \mathcal{G}_d .

By definition the dimension n of a geometry is the maximal number of linear independent vectors. In the given case dimension of \mathcal{G}_E is equal to 2, as well as the dimension of the manifold, where the geometry is given. Dimension of the manifold is defined as the number of coordinates of the coordinate system.

Dimension of the manifold in the physical geometry \mathcal{G}_d is also 2, as well as in \mathcal{G}_E . In the physical geometry (T-geometry) m vectors $\mathbf{P}_0\mathbf{P}_1, \mathbf{P}_0\mathbf{P}_2, \dots, \mathbf{P}_0\mathbf{P}_m$ are linear independent if and only if the Gram's determinant

$$F_m(\mathcal{P}^m) \neq 0, \quad F_m(\mathcal{P}^m) \equiv \det \|(\mathbf{P}_0\mathbf{P}_i, \mathbf{P}_0\mathbf{P}_k)\|, \quad i, k = 1, 2, \dots, m \quad (3.6)$$

Here $\mathcal{P}^m = \{P_0, P_1, \dots, P_m\}$, and the scalar product $(\mathbf{P}_0\mathbf{P}_i, \mathbf{Q}_0\mathbf{Q}_k)$ of two vectors $\mathbf{P}_0\mathbf{P}_1, \mathbf{Q}_0\mathbf{Q}_1$ is defined by the relation (1.3).

The conventional definition of linear independence is different. s vectors $\mathbf{P}_0\mathbf{P}_1, \mathbf{P}_0\mathbf{P}_2, \dots, \mathbf{P}_0\mathbf{P}_s$ are linear independent, if the linear combination of vectors satisfies the relation

$$\sum_{k=1}^{k=s} \alpha_k \mathbf{P}_0\mathbf{P}_k = 0 \quad (3.7)$$

only at $\alpha_k = 0, \quad k = 1, 2, \dots, s$. For the proper Euclidean geometry both definitions (3.6) and (3.7) are equivalent. For the distorted geometry \mathcal{G}_d they are not equivalent, in general.

The conventional definition (3.7) supposes existence of the linear vector space with a scalar product, given on it, and, in particular, it supposes the procedures of definition : summation of vectors and multiplication of a vector by a real number. Definition (3.6) contains references only to the world function and points of the space. Existence of the linear vector space and linear operations over vectors is not supposed. It is evident, that the definition (3.6) is a more general definition, than

(3.7), which can be applied, only if the linear vector space can be introduced. It seems rather unexpected, that one can define linear dependence, without introduction of the linear space, because the name "linear dependence" implicates conventionally existence of the linear space. However, the definition (3.6) can be used in the case, when one cannot introduce the linear space. In this case the determinant, constructed of scalar products of vectors, describes interrelations of m vectors, in particular, their mutual orientation.

Let us consider four vectors

$$\mathbf{P}_0\mathbf{P}_1 = \{a, 0\}, \quad \mathbf{P}_0\mathbf{P}_2 = \{0, b\}, \quad \mathbf{P}_0\mathbf{P}_3 = \{a, b\}, \quad \mathbf{P}_0\mathbf{P}_2 = \{2a, 0\} \quad (3.8)$$

Let us suppose for simplicity, that coordinates $a, b \gg \lambda_0$. Then the scalar products of any vectors (3.8) in the proper Euclidean geometry \mathcal{G}_E and in the distorted geometry \mathcal{G}_d are connected by the relations

$$(\mathbf{P}_0\mathbf{P}_i, \mathbf{P}_0\mathbf{P}_k)_d = (\mathbf{P}_0\mathbf{P}_i, \mathbf{P}_0\mathbf{P}_k)_E - 2\lambda_0^2, \quad \text{if } P_i \neq P_k \quad (3.9)$$

$$(\mathbf{P}_0\mathbf{P}_i, \mathbf{P}_0\mathbf{P}_i)_d = (\mathbf{P}_0\mathbf{P}_i, \mathbf{P}_0\mathbf{P}_i)_E - \lambda_0^2 \quad (3.10)$$

These relations are the corollaries of the relations (3.4) and (1.3)

The Gram's determinant in \mathcal{G}_d for the first three vectors (3.8) has the form

$$\begin{vmatrix} a^2 - 2\lambda_0^2 & -\lambda_0^2 & a^2 - \lambda_0^2 \\ -\lambda_0^2 & b^2 - 2\lambda_0^2 & b^2 - \lambda_0^2 \\ a^2 - \lambda_0^2 & b^2 - \lambda_0^2 & a^2 + b^2 - 2\lambda_0^2 \end{vmatrix} = -4\lambda_0^2 (b^2 - \lambda_0^2) (a^2 - \lambda_0^2) \quad (3.11)$$

For the four vectors (3.8) the Gram's determinant in \mathcal{G}_d has the form

$$\begin{vmatrix} a^2 - 2\lambda_0^2 & -\lambda_0^2 & a^2 - \lambda_0^2 & 2a^2 - \lambda_0^2 \\ -\lambda_0^2 & b^2 - 2\lambda_0^2 & b^2 - \lambda_0^2 & b^2 - \lambda_0^2 \\ a^2 - \lambda_0^2 & b - \lambda_0^2 & a^2 + b^2 - 2\lambda_0^2 & 2a^2 + b^2 - \lambda_0^2 \\ 2a^2 - \lambda_0^2 & b^2 - \lambda_0^2 & 2a^2 + b^2 - \lambda_0^2 & a^2 + b^2 - 2\lambda_0^2 \end{vmatrix} \\ = -\lambda_0^2 (12a^4\lambda_0^2 - 12a^4b^2 + 2a^2\lambda_0^4 + 6b^2\lambda_0^4 - 5\lambda_0^6 - 3a^2b^2\lambda_0^2) \quad (3.12)$$

In \mathcal{G}_d the Gram's determinant for two "collinear" vectors $\mathbf{P}_0\mathbf{P}_1 = \{a, 0\}$, $\mathbf{P}_0\mathbf{P}_2 = \{2a, 0\}$ has the form.

$$\begin{vmatrix} a^2 - \lambda_0^2 & 2a^2 - 2\lambda_0^2 \\ 2a^2 - 2\lambda_0^2 & 4a^2 - \lambda_0^2 \end{vmatrix} = 3\lambda_0^2 (a^2 - \lambda_0^2) \quad (3.13)$$

although in the Euclidean geometry \mathcal{G}_E this determinant vanishes. In general, in the geometry \mathcal{G}_E all three determinants (3.11), (3.12), (3.13) vanish, because $\lambda_0^2 = 0$ and dimension of the geometry \mathcal{G}_E is equal to 2.

It follows from (3.11) and (3.12), that in the distorted geometry \mathcal{G}_d there are, at least, four linear independent vectors, although the dimension of the manifold remains to be equal to 2. One should expect, that in the distorted geometry \mathcal{G}_d there is infinite number of linear independent vectors, and concept of dimension is inadequate for the physical geometry \mathcal{G}_d .

Thus, in the proper Euclidean geometry the dimension $n_{\mathcal{M}}$ of a manifold is equal to the dimension $n_{\mathcal{G}}$ of the geometry, whereas in the distorted geometry \mathcal{G}_d dimension of the manifold and dimension of the geometry are quite different quantities. It looks rather unexpected. How is it possible?

The dimension $n_{\mathcal{G}}$ of a geometry is a very complicated concept, but it concerns to the geometry itself. The dimension $n_{\mathcal{M}}$ of a manifold is a simple concept, but it relates only to the method of description (manifold). In the proper Euclidean geometry the values (but not concepts) of the two dimensions coincide ($n_{\mathcal{G}} = n_{\mathcal{M}}$). Conventionally one does not distinguish between the two dimensions. It leads to a confusion and to an ascription of the description properties to the geometry in itself.

The dimension $n_{\mathcal{M}}$ of manifold may be defined only for a continuous set of space points. It is invariant only with respect to continuous coordinate transformation. In this connection it is interesting a consideration of the discrete geometry \mathcal{G}_{dis} . Let us consider the two-dimensional proper Euclidean geometry, given on the point set Ω . The point set Ω_2 is obtained from the point set Ω as follows. Let K_2 be a Cartesian coordinate system on Ω . Let us remove all points of Ω , except of those points, for which both Cartesian coordinates are integer. The remaining point set Ω_2 forms a grid. World function σ_{dis} is defined on the set $(\Omega_2 \times \Omega_2) \subset (\Omega \times \Omega)$. On this set the world function σ_{dis} coincides with σ_{E} and, hence, it satisfies to all conditions of Euclideaness [1] except for the condition IV (the continuity condition). Dimension $n_{\mathcal{G}}$ of geometry \mathcal{G}_{dis} , determined by means of the definition (3.6), is equal to 2. Dimension $n_{\mathcal{M}}$ of the "manifold" Ω_2 cannot be determined definitely, because the number of integer variables, labeling the points of Ω_2 may be 1, 2, ... The dimension $n_{\mathcal{M}}$ is invariant only with respect to continuous coordinate transformation. In the case, when coordinates are integer, there are no continuous coordinate transformations. In this case the dimension of manifold has no sense, because there is no manifold, whereas the geometry dimension $n_{\mathcal{G}}$ is defined correctly in the case of the discrete geometry.

4 Multivariate, discreteness and graininess of space-time

Conventionally the discrete geometry is considered on some grid of points. It seems that a geometry, given on a continuous manifold, cannot be discrete. It means that conventionally a discreteness of a geometry is connected with the means of the geometry description, (but not with the geometry in itself). In reality, the discreteness of the geometry is determined by the world function. In particular, a discrete geometry may be given on the continuous manifold. Besides, there may be different degrees of the physical geometry discreteness.

Let us consider the question on discreteness of the space-time geometry, described by the world function

$$\sigma_d = \sigma_M + d \cdot \text{sgn}(\sigma_M), \quad d = \lambda_0^2 = \text{const} > 0 \quad (4.1)$$

$$\text{sgn}(x) = \begin{cases} 1, & \text{if } x > 0 \\ 0, & \text{if } x = 0 \\ -1, & \text{if } x < 0 \end{cases}, \quad (4.2)$$

where σ_M is the world function of the 4-dimensional space-time of Minkowski. λ_0 is some elementary length. The space-time geometry (4.1) is closer to the real space-time geometry of microcosm, than the space-time of Minkowski, because at this space-time geometry the quantum effects may be described without a use of the quantum principles, if the elementary length $\lambda_0 = \hbar^{1/2} (2bc)^{-1/2}$. Here c is the speed of the light, \hbar is the quantum constant, and b is some universal constant, whose exact value is not determined [4].

The space-time geometry (4.1) is a discrete space-time geometry, because in this space-time geometry there are no vector $\mathbf{P}_0\mathbf{P}_1$, whose length $|\mathbf{P}_0\mathbf{P}_1|$ be small enough, i.e.

$$|\mathbf{P}_0\mathbf{P}_1|^4 \notin (0, \lambda_0^4), \quad \forall P_0, P_1 \subset \Omega \quad (4.3)$$

In other words, the space-time geometry (4.1) has no close points.

Let us consider another space-time geometry \mathcal{G}_d which is partly discrete. World function σ_d of this geometry has the form

$$\sigma_d = \sigma_M + d(\sigma_M) \quad (4.4)$$

$$d(\sigma_M) = \lambda_0^2 f\left(\frac{\sigma_M}{\sigma_0}\right) = \begin{cases} \lambda_0^2 \text{sgn}\left(\frac{\sigma_M}{\sigma_0}\right) & \text{if } |\sigma_M| > \sigma_0 > 0 \\ \lambda_0^2 \frac{\sigma_M}{\sigma_0} & \text{if } |\sigma_M| \leq \sigma_0 \end{cases} \quad (4.5)$$

where σ_M is the world function of the geometry of Minkowski.

If σ_0 is small, the world function is close to the world function (4.1). If $\sigma_0 \rightarrow 0$, the world function (4.5) tends to (4.1). Strictly, the space-time geometry (4.5) is not discrete, however it is close to the discrete space-time geometry (4.1).

Let us consider the relative density $\rho(\sigma_d) = d\sigma_d/d\sigma_E$ of points of the geometry \mathcal{G}_d with respect to the geometry \mathcal{G}_E . One obtains

$$\rho(\sigma_d) = d\sigma_d/d\sigma_E = \begin{cases} 1 & \text{if } |\sigma_d| > \sigma_0 + \lambda_0^2 \\ \frac{\sigma_0}{\sigma_0 + \lambda_0^2} & \text{if } |\sigma_d| \leq \sigma_0 + \lambda_0^2 \end{cases} \quad (4.6)$$

If $\sigma_0 = 0$, there is no close points which are separated by interval with the world function $\sigma_d \in (0, \lambda_0^2)$ and $\sigma_d \in (-\lambda_0^2, 0)$. It means, that the space-time geometry is discrete at $\sigma_0 = 0$.

If $\sigma_0 \ll \lambda_0^2$, the relative density $\rho(\sigma_d) \simeq \sigma_0/\lambda_0^2$ of points inside the interval $\sigma_d \in (-\sigma_0 - \lambda_0^2, \sigma_0 + \lambda_0^2)$ is much less, than unity. It means that space-time geometry is almost discrete. The quantity $1 - \rho(\sigma_d)$, $\sigma_d \in (-\sigma_0 - \lambda_0^2, \sigma_0 + \lambda_0^2)$ may be interpreted as a degree of the discreteness of the space-time geometry. One can see, that the discreteness of the space-time geometry and the degree of the discreteness is determined by properties of the world function (but not by properties of the manifold). The fact, that the space-time geometry, given on a continuous manifold may be discrete, seems to be very unexpected. This fact acknowledges the statement, that the world function and only world function determines the space-time geometry.

It is reasonable to interpret the relative density $\rho(\sigma_d) = d\sigma_d/d\sigma_E$ of points of the distorted space-time with respect to the density of points of the standard (Minkowskian) space-time as a measure of the space-time granulation. Discreteness is a special case of graininess. Continuity is another special case of graininess. The graininess of the space-time describes also all intermediate cases, when the space-time is partly continuous and partly discrete. Interrelation of the graininess with the discreteness reminds interrelation of rational numbers with the natural ones.

Graininess of the space-time is a physical property of the space-time, whereas the multivariance is a mathematical property of the space-time geometry. The graininess of the space-time is connected with the multivariance, and the world function formalism is a mathematical technique, which can describe the graininess of the space-time.

One can easily imagine two limit cases of graininess: discreteness and continuity. The relative density $\rho(\sigma_d) = d\sigma_d/d\sigma_E$ admits one to realize intermediate cases of graininess. Conventional approach to geometry, based on the concept of linear vector space, describes only continuous geometries. Vector representation of geometry [3], based on the concept of the linear vector space, cannot describe indefinite graininess of the space-time. No finesses, based on the vector representation of geometry, enables to describe effectively the space-time, whose graininess distinguishes from continuity.

The discrete values of the elementary particles characteristics (mass, charge, spin) are generated by some discrimination mechanism. The reason of this discrimination is conditioned by the multivariance (more exactly by zero-variance) of the space-time geometry [13]. From physical viewpoint the reason of the discrete characteristics is the graininess of the space-time.

5 Concluding remarks

Thus, the multivariance is a general property of the space-time geometry. Class of uniform, isotropic flat space-time geometries is a continual set, whose elements are labelled by a real function of one argument. Only one geometry of this class may be considered as single-variant (geometry of Minkowski). All other space-time geometries are multivariant. Considering the Riemannian geometry as a more general space-time geometry, we restrict our capacities. At the conventional approach, based on concepts of the linear vector space, the natural multivariance of the Riemannian geometry is suppressed by means of the fernparallelism interdict.

In the framework of the Riemannian geometry one cannot describe such properties of the space-time as the limited divisibility and the graininess, which are very important for geometrical description of elementary particles. In general, ignoring multivariant geometries, we manifest, that our knowledge of geometry are very restrictive. Our knowledge of geometry does not admit one to construct effective description and effective dynamics of elementary particles in microcosm, where the graininess of the space-time is important.

The multivariance and the graininess are connected between themselves. However, the graininess is rather physical concept, whereas the multivariance is rather mathematical concept. The multivariance describes the interrelation of two vectors, whereas the graininess describes interrelation of the point density in the space-time with the standard point density in the space-time of Minkowski. The graininess is more demonstrative and more complicate, whereas the multivariance is less demonstrative and simpler. As a result the multivariance is considered as a basic concept of the space-time, whereas the graininess is considered as a derivative concept.

In the Riemannian geometry the unlimited divisibility of the space-time takes place. As a result, on one hand, the particle dynamics can be described in terms of differential equations, whose application supposes the unlimited divisibility of the space-time. On the other hand, the unlimited divisibility generates such problems as confinement

One can hardly formulate mathematically the particle dynamics in the grainy space-time, where the space-time divisibility is restricted, and one cannot use differential equations. In the grainy space-time the particle dynamics is determined by the space-time geometry in itself and by the structure of the particle. Such a geometric dynamics is formulated in terms of the world chain with finite links [14]. The world chain is such a generalization of the world line, when infinitesimal segments of the world line are replaced by finite geometrical objects.

References

- [1] Geometry without topology as a new conception of geometry. *Int. Journ. Mat. & Mat. Sci.* **30**, iss. 12, 733-760, (2002), (See also *e-print* <http://arXiv.org/abs/math.MG/0103002>).
- [2] D. Hilbert, *Grundlagen der Geometrie*. 7 Auflage, B.G.Teubner, Leipzig, Berlin, 1930.
- [3] Yu.A.Rylov, Different conceptions of Euclidean geometry. *e-print* <http://arXiv.org/abs/0709.2755>
- [4] Yu. A. Rylov, Non-Riemannian model of space-time, responsible for quantum effects, *J. Math. Phys.* **32** (8), 2092 - 2098, (1991)
- [5] Yu. A. Rylov, Tubular geometry construction as a reason for new revision of the space-time conception. (Printed in *What is Geometry?* polimetrica Publisher, Italy, pp.201-235).
- [6] Yu. A.Rylov, Non-Euclidean method of the generalized geometry construction and its application to space-time geometry in *Pure and Applied Differential geometry PADGE 2007*, pp.238-246. eds. Franki Dillen and Ignace Van de Woestyne. Shaker Verlag, Aachen, 2007. (or *e-print* <http://arXiv.org/abs/Math.GM/0702552>)

- [7] Yu. A. Rylov, New crisis in geometry? e-print <http://arXiv.org/abs/math.GM/0503261>.
- [8] G. Perelman, The entropy formula for the Ricci flow and its geometric applications. *e-print* <http://arXiv.org/abs/math.DG/0211159>
- [9] G. Perelman, Ricci flow with surgery on three-manifolds. *e-print* <http://arXiv.org/abs/math.DG/0303109>.
- [10] G. Perelman, Finite extinction time for the solutions to the Ricci flow on certain three-manifolds. *e-print* <http://arXiv.org/abs/math.DG/0307245>.
- [11] A. Abramov, *Moscow News*, Number 32, 2006.
- [12] Yu. A. Rylov, Euclidean geometry as algorithm for construction of generalized geometries. *e-print* <http://arXiv.org/abs/math.GM/0511575>.
- [13] Yu. A. Rylov, Discrimination of particle masses in multivariant space-time geometry *e-print* <http://arXiv.org/abs/0712.1335>.
- [14] Yu. A. Rylov, Geometrical dynamics: spin as a result of rotation with superluminal speed. *e-print* <http://arXiv.org/abs/0801.1913>.